HIGH-PRECISION ℓN COLLISIONS: BASICS OF RADIATIVE CORRECTIONS

EINN 2009

Workshop on Electron-Ion Colliders

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MOTIVATION

The Goal:

High-Precision Measurements of the 'Nucleon Structure'

➜ measure form factors, structure functions, (generalized) parton distribution functions, ...

- at low Q^2 elastic and quasi-elastic scattering
  ➜ form factors, polarizabilities, ...
- at high Q^2 deep inelastic scattering
  ➜ parton distribution functions, GPDs, GDAs, ...

The interesting physics is encoded in FFs, PDFs, ...
test the dynamics of the strong interaction

Precision Physics (QCD) D. de Florian

Lepton scattering: only via electromagnetic and weak interaction

➜ well-controlled and separable perturbative treatment

RADIATIVE CORRECTIONS

Measure FFs, PDFs, etc by comparing data with theoretical predictions:

$$\sigma_{\text{exp}} = \sigma_{\text{theory}}[F_n(x, Q^2, \ldots)]$$

High precision requires knowledge of higher-order corrections

$$\sigma_{\text{theory}} = \sigma^{(0)} + \alpha_{\text{em}}\sigma^{(1)} + \ldots$$

- Emission of real photons
  often not distinguished from non-radiative processes:
  soft photons, collinear photons
  ➜ "radiative corrections"
- Virtual corrections (loop diagrams)
  needed to cancel infrared divergences (Bloch-Nordsieck)
- Electroweak effects
  Z-, W-boson exchange ($O(G_F)$)
  and higher-order electroweak corrections ($O(\alpha G_F)$)

Radiative corrections have to be 'removed' to uncover the interesting physics and radiative corrections often considered the uninteresting part

but:

- 2-γ-exchange (and γZ-exchange) is part of the 1-loop photonic corrections
  corresponding infrared divergences cancel with the interference between photon radiation from the lepton and from the nucleon
- radiation from the nucleon: DVCS deeply virtual Compton scattering, is part of the 1-photon radiative corrections
- Z-exchange gives rise to P- and C-violating interactions, charge and polarization asymmetries

CLASSIFICATION OF $O(\alpha)$ CORRECTIONS

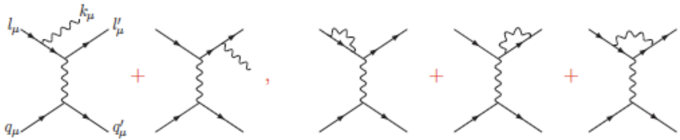
- Radiation from the lepton
  model independent (universal)
- vacuum polarization (boson self energy)
  universal
- Radiation from the nucleon
  parton model: radiation from quarks
  to be considered as a part of the nucleon structure
- Interference of leptonic and hadronic radiation
  2γ exchange
  new structure
- purely weak corrections

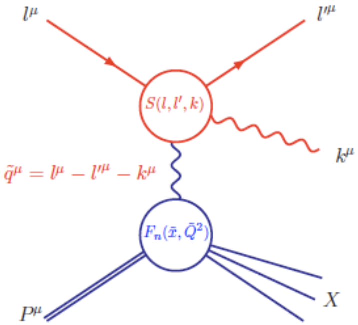
Note: for NC-scattering straightforward separation
Rule: respect gauge invariance
IR divergences: need to combine real and virtual radiation

Leptonic radiation

Feynman diagrams for leptonic radiation at $O(\alpha)$

for *eq* scattering:





radiative leptonic tensor $S_{\mu\nu}(l, l', k)$ is

- gauge invariant
- infrared finite
- universal

Observed cross section: convolution of true cross section $\otimes$ radiator function

$$\mathrm{d}\sigma^{\mathrm{obs}}(p,q) = \int \frac{d^3k}{2k^0} R(l,l',k) \mathrm{d}\sigma^{\mathrm{true}}(p,q-k)$$

of, for the structure functions:

$$F_n^{\mathrm{obs}}(x,Q^2) = \int \mathrm{d}\tilde{x}\mathrm{d}\tilde{Q}^2 R_n(x,Q^2;\tilde{x},\tilde{Q}^2) F_n^{\mathrm{true}}(\tilde{x},\tilde{Q}^2)$$

Note: shifted kinematics, e.g.,

$$Q^2 = -(l-l')^2 \rightarrow \tilde{Q}^2 = -(l-l'-k)^2$$

➜ expect strong dependence on experimental prescriptions for measuring kinematic variables (➜ no numerical results in this talk)
in turn: measurement = unfolding may be ill-defined
Difficult to treat radiative and detector effects separately (Acceptance cuts, efficiencies, ...)

with partial fractioning, write: $R(l, l', k) = \frac{I}{k \cdot l} + \frac{F}{k \cdot l'}$

- initial state radiation, $k \cdot l$ small for $\sphericalangle(e_{\rm in}, \gamma) \to 0$
- final state radiation, $k \cdot l'$ small for $\sphericalangle(e_{\rm out}, \gamma) \to 0$

narrow peaks, width $\simeq \sqrt{m_e/E_e}$: collinear or mass singularities

upon angular integration: large logarithm $\propto \frac{\alpha}{\pi} \log \frac{Q^2}{m_e^2} \simeq 10\,\%$

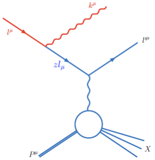
Note: $E^2_{\gamma,\rm max} \propto Q^2 \frac{1-x}{x}$

➜ large corrections at large Q^2 and at small x

➜ Radiation suppressed at small Q^2 and at large x,
large negative corrections from uncancelled virtual contributions

e.g., for initial-state radiation:
assume $k^\mu = (1-z)l^\mu$
➜ Radiator function

$$R_{\rm ISR} = \frac{\alpha}{2\pi}\,\frac{1+z^2}{1-z}\,\log\frac{Q^2}{m_e^2}$$



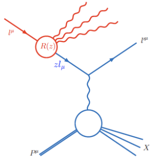
$(+\delta(1-z)$ from loops $\rightarrow$ $+$-distribution $1/(1-z)_+)$

$$d\sigma_{\rm ISR} = \int \frac{dz}{z} R_{\rm ISR}(z)\, d\sigma_{\rm Born}(l^\mu \rightarrow z l^\mu)$$

(similar for final-state radiation)

Can be extended to include multi-photon emission:

$$R^{(2)}_{\rm ISR}(z) = \int_z^1 \frac{dz'}{z'} R^{(1)}_{\rm ISR}(z') R^{(1)}_{\rm ISR}(z/z') + \ldots$$



Solution of evolution equations like DGLAP
Known at $O(\alpha^2)$ (complete) and partially at $O(\alpha^3)$

Exponentiation

Corrections due to soft photons are universal

sum of real and virtual contributions: $\delta^{\rm IR}$ (finite and gauge invariant)

$$1 + \delta^{\rm tot} = 1 + \delta^{\rm IR} + \delta^{\rm fin} \rightarrow \exp(\delta^{\rm IR})(1 + \delta^{\rm fin})$$

$\delta^{\rm IR}$ contains $\log(E_\gamma^{\rm max})$ and $L_f = \log(m_f^2/Q^2)$:

$$1 + \frac{\alpha}{2\pi}(L_f - 1)\ln\frac{E_\gamma^{\rm max}}{E_e} + \ldots \rightarrow \left(\frac{E_\gamma^{\rm max}}{E_e}\right)^{\frac{\alpha}{2\pi}(L_f-1)} (1 + \ldots)$$

(in the $\gamma^* p$ cms: $E_\gamma^{\rm max} = \frac{1}{2}\sqrt{y(1-x)S}$, i.e. important at low y and large x)

Yennie, Frautschi, Suura, 1961

Radiative tail

Radiation of (hard) photons ➜ shifted kinematic variables:

$$Q^2 = -(l - l')^2 \rightarrow \tilde{Q}^2 = -(l - l' - k)^2$$

and

$$x = \frac{Q^2}{2P \cdot (l - l')} \rightarrow \tilde{x} = \frac{\tilde{Q}^2}{2P \cdot (l - l' - k)}$$

Radiator function is folded with

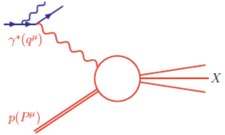
$$d\sigma(\tilde{x}, \tilde{Q}^2) \propto \frac{1}{\tilde{Q}^2}$$

➜ correction factor $d\sigma_{O(\alpha)}(x, Q^2)/d\sigma_{\text{Born}}(x, Q^2)$ enhanced by $Q^2/\tilde{Q}^2$

Note: $\tilde{Q}^2 \ll Q^2$ possible: $\tilde{Q}^2_{\min} = \frac{x^2}{1-x} M_N^2$

➜ radiative tail, Compton peak

back to photoproduction



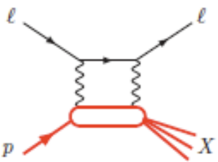
TWO-PHOTON EXCHANGE

(deep) inelastic *ep*:

factorized into

PDFs $\otimes$ 2γ-box

for *eq* scattering



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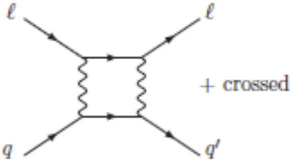
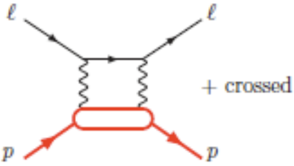


Need intereference of radiation from the lepton and the hadron to obtain IR-finite result

elastic *ep*:

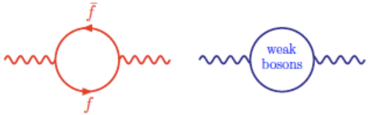
- assume dominance of a few intermediate states: *p* + resonances

- assume factorization into GPDs $\otimes$ partonic scattering



Dedicated precision measurements to determine TPE:
lepton charge asymmetry (*Re*) and lepton polarization asymmetry (*Im*)
see e.g. M.Kohl, K. de Jager, MAMI&Beyond

Self energy diagrams of the exchanged boson (γ and Z)



$\propto \log \frac{Q^2}{m_f^2} \to O(10\,\%)$ small, $O(1\,\%)$

Photon self energy = vacuum polarization, absorbed in the running fine structure constant:

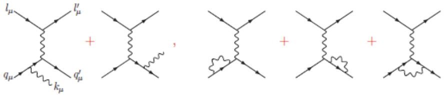
$$\alpha \to \alpha(Q^2) = \frac{\alpha}{1 - \Pi_\gamma(Q^2)}$$

Z-boson self energy: a small correction if written in terms of:

$$\frac{\alpha}{s_W^2 c_W^2} \to \frac{M_Z^2 G_\mu \sqrt{2}}{\pi} \frac{1 - \Delta r}{1 - \Pi_Z(Q^2)}$$

(with s_W^2, c_W^2: sin and cos of the weak mixing angle; G_μ the muon decay constant; Δr one-loop corrections to the muon decay: renormalization)

at large Q^2: DIS, parton model
emission of photons
like emission of gluons



infrared divergences (soft photons / gluons) cancel with loops, collinear emission gives rise to corrections $\frac{\alpha}{2\pi} \log m_q^2$, but quark masses are ill-defined

➜ factorize and absorb collinear divergences into parton distribution functions

$$d\sigma = \sum_f d\hat{\sigma}_f (1 + \delta_f(Q^2; m_q^2)) q_f(x)$$

$$d\sigma = \sum_f d\hat{\sigma}_f (1 + \delta_f(Q^2; m_q^2)) q_f(x) = \sum_f d\hat{\sigma}_f \hat{q}_f(x, Q^2)$$

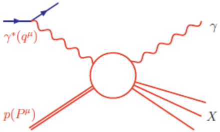
renormalized parton distribution functions

$$\hat{q}_f(x, Q^2) = (1 + \delta_f(Q^2; m_q^2)) q_f(x)$$

➜ modified scaling violations

well-known in QCD, $\overline{\text{MS}}$ factorization

Different point of view:
not as part of radiative corrections
but with observed photon



➜ DVCS: deeply virtual Compton scattering

see e.g. B. Pasquini, MAMI&Beyond

MONTE-CARLO IMPLEMENTATION

classical analytical approach: Mo, Tsai
often used in 'private' implementations of experimental collaborations

full Monte-Carlo approach:
HERACLES: complete electroweak corrections at $O(\alpha)$ (parton model) for NC and CC scattering at HERA, including polarization

full event generation:
DJANGO: universal leptonic corrections at $O(\alpha)$, interface to QCD-based event generation of jets, parton showers, hadronic final state, includes models for low Q^2 behaviour: elastic tail, SOPHIA for low-mass hadronic final states

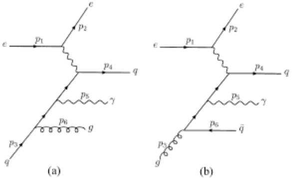
specialized: VANDERHAEGHEN ET AL. QED $O(\alpha)$ corrections to virtual Compton scattering

... and many more tools: TERAD by Bardin et al.,
HEKTOR, KRONOS, FRANEQ, ...

MONTE-CARLO: FUTURE IMPROVEMENTS ?

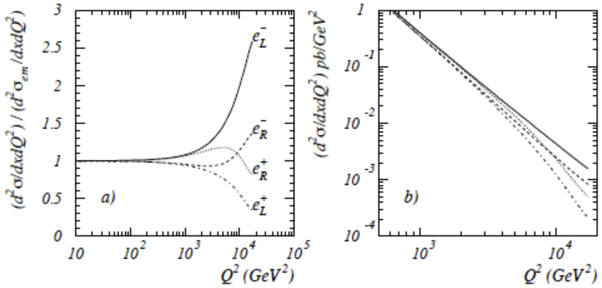
More dedicated efforts needed to include:

- IR/soft photon exponentiation
- multi-photon emission
  radiator functions at $O(\alpha^2)$
- 2-photon exchange
- radiation from quarks:
  subtraction and modified parton showers
  including $q \to q + \gamma$
  (mixed QED+QCD corrections)



Z-boson exchange

Charge and polarization asymmetries see e.g. J. Erler, MAMI&Beyond



$$A_\pm = \frac{\sigma(e_L^\pm) - \sigma(e_R^\pm)}{\sigma(e_L^\pm) + \sigma(e_R^\pm)}, \quad B_\pm = \frac{\sigma(e_L^\pm) - \sigma(e_R^\mp)}{\sigma(e_L^\pm) + \sigma(e_R^\mp)}, \quad C_{L,R} = \frac{\sigma(e_{L,R}^-) - \sigma(e_{L,R}^+)}{\sigma(e_{L,R}^-) + \sigma(e_{L,R}^+)}$$

Full $O(\alpha)$ electroweak corrections based on a parton model calculation available

SUMMARY

- High precision needs careful treatment of radiative corrections
- Closely related to experimental conditions
- Interesting physics: DVCS, TPE, electroweak effects